

### §31. Theory of Anomalous Transport in the Reverse Field Pinch

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Recently, efforts have been made to understand the anomalous transport coefficient in toroidal plasmas. A new method of plasma turbulence and anomalous transport, i.e., the self-sustained turbulence, seems to provide agreement with observations of various features [1]. It is also meaningful to examine the applicability of the transport theory to the reverse field pinch.

We employ the model of the high aspect ratio RFP. The toroidicity is neglected and the magnetic field is given by the Bessel function model as

$$B_t = B_0 J_0(\zeta r), B_p = B_0 J_1(\zeta r) .$$

The magnetic field is characterized by parameters  $B_0$  and  $\zeta$ . The parameter  $\zeta$  is in the range of  $\zeta a \approx 3$ . ( $a$  being the minor radius.) This configuration has the magnetic hill. The magnetic curvature  $\kappa$  and shear  $s$  are calculated as

$$\kappa = \frac{a}{B} \frac{dB}{dr} \approx -\frac{\zeta^2 a}{4}$$

$$s \approx \frac{Rr}{B_t} \frac{d}{dr} \left( \frac{B_p}{r} \right) \approx -\frac{1}{8} \zeta^3 R r^2$$

near the origin ( $R$  being the major radius.). Near the field reversal surface, we have  $s \approx -\zeta R$ , and  $s$  remains of the order of  $R/a$ .

In the presence of magnetic braiding, the thermal transport coefficient of electrons was obtained as

$$\chi = \frac{1}{s^2} |R\beta' \kappa R/a|^{3/2} \left( \frac{c}{\omega_p} \right)^2 \frac{v_A}{R} M$$

where  $\beta$  is the ratio of the plasma pressure to the total magnetic pressure,  $v_A$  is the Alfvén velocity and  $\omega_p$  is the electron plasma frequency. The coefficient  $M$  denotes the enhancement factor of the magnetic braiding in

evaluating  $\chi_e$ . In the electrostatic limit, we have  $M = 1$ , and  $M \approx m_i/3m_e$  for the case of magnetic stochasticity [2]. The condition for magnetic braiding has been obtained as  $\beta \geq \beta_c$  with the critical pressure gradient  $\beta_c \approx 4 / \zeta^2 a R$ . The condition for magnetic braiding is satisfied in the range of interest,  $\beta_p \approx 0.2$ , found in experimental observations.

The energy balance equation for the Ohmic heating plasma is given as

$$\langle \chi \rangle a^{-2} \langle n \rangle \langle T \rangle \approx \langle \eta \rangle \langle J \rangle^2$$

where  $\eta$  is the resistivity and  $J$  is the current density. We write the resistivity as  $\eta = \eta_0 Z_{\text{eff}} T^{-3/2}$  where  $\eta_0$  is a universal constant and  $Z_{\text{eff}}$  includes both the effects of impurity collisions and a possible anomalous resistivity due to plasma turbulence.

From the energy balance equation, we have

$$T \propto Z_{\text{eff}}^{1/4} a^{-3/4} I_p n^{-1/4}$$

$$\beta_p \propto M^{-1/4} Z_{\text{eff}}^{1/4} a^{-1/4} N^{3/4} I_p^{-1}$$

where  $N$  is the number per unit length,  $N = na^2$  and  $I_p$  is the plasma current.

These results are compared with experimental observations[3]

$$\beta_p \propto N / I_p$$

$$T \propto I_p^{0.8} n^{-0.2}$$

which agree well with the theoretical prediction.

#### References

- 1) Fukuyama A, Itoh K, Itoh S-I, Yagi M, Azumi M 1994 *Plasma Phys. Contr. Fusion* **36** 1385.
- 2) Itoh K, Fukuyama A, Itoh S-I, Yagi M, Azumi M 1994 "Self-sustained magnetic braiding in toroidal plasmas" Research Report NIFS-288.
- 3) Fujiwara M 1994 in *Plasma Physics and Controlled Nuclear Fusion Research 1992* Vol.4 p.23.